

Discrete Fourier Transforms

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Fourier transform

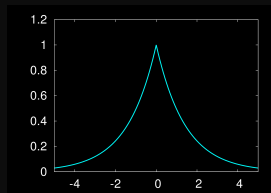
In this lecture, we will discuss:

- The Fourier transform,
- The discrete Fourier transform
- The **fast** Fourier transform
- Examples using the FFTW library



Fourier transform

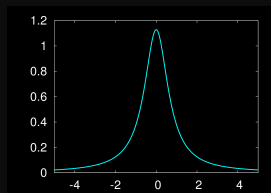
- Let f be a function of a spatial variable x .



$$f(x) = e^{-|x|}$$

- Transform to a function \hat{f} of the wavenumber k :

$$\hat{f}(k) \propto \int f(x) e^{\pm i k x} dx$$



$$f(k) = (1 + k^2)^{-1}$$

- Inverse transformation:

$$f(x) \propto \int \hat{f}(k) e^{\mp i k x} dk$$

Fourier transform

- Fourier made the claim that any function can be expressed as a harmonic series.
- The Fourier Transform is a mathematical expression of that.
- Constitutes a linear basis transformation in function space.
- Transforms from spatial to wavenumber, or time to frequency, etc.
- Constants and signs are just convention.*

* some restrictions apply.

Finite interval = periodicity

- Imagine the function is stored numerically as n samples at $x = j\Delta x$.
- Then the function lives on a **finite** x -interval $[0, L)$ where $L = n\Delta x$.
- For a finite interval, only discrete k values are needed, e.g.

$$\hat{f}(k) \propto \int f(x) e^{\pm i k x} dx \equiv \hat{f}_q \quad \text{for } k = \frac{2\pi}{L} q$$

- This is enough to reconstruct the inverse

$$f(x) \propto \sum_{q=-\infty}^{\infty} \hat{f}(k(q)) e^{\mp i k(q) x} dx$$

- $k(q)$ is such that $e^{\pm i k(q) x}$ has period L ,
so we can view the function f as **periodic** in x with period L .

Discrete case

- We only have the n discrete points $j\Delta x$.
- The Fourier transform is invertible, so it must have the same number of points.
- We'll keep just $q = 0$ up to $q = n - 1$. Then:

$$f_j \equiv f(j\Delta x) \propto \sum_{q=0}^{n-1} \hat{f}(k(q)) e^{\mp i k(q) j \Delta x} dx = \sum_{q=0}^{n-1} \hat{f}_q e^{\mp 2\pi i q j} dx$$

- Higher q values coincide with lower fourier modes as far as $x = j\Delta x$ is concerned.
- But we have to alter how we compute \hat{f} ; the integral must become a sum:

$$\hat{f}_q = \sum_{j=0}^{n-1} f_j e^{\pm i j \Delta x k(q)} = \sum_{j=0}^{n-1} f_j e^{\pm 2\pi i j q/n}$$

Discrete Fourier Transform (DFT)



C. F. Gauss

- Given a set of n function values on a regular grid:

$$x_j = j\Delta x; \quad f_j = f(j\Delta x)$$

- Transform to n other values

$$\hat{f}_q = \sum_{j=0}^{n-1} f_j e^{\pm 2\pi i j q/n}$$

- Easily back-transformed:

$$f_j = \frac{1}{n} \sum_{q=0}^{n-1} \hat{f}_q e^{\mp 2\pi i j q/n}$$

- Note: $\hat{f}_{-q} = \hat{f}_{n-q}$.
- Note: Cannot resolve frequencies higher than the Nyquist frequency $q = n/2$ ($k = \pi/\Delta x$).

Slow Fourier transform

$$\hat{f}_q = \sum_{j=0}^{n-1} f_j e^{\pm 2\pi i j q/n}$$

- Discrete fourier transform is a linear transformation.
- In particular, it's a matrix-vector multiplication.
- Naively, costs $\mathcal{O}(n^2)$. Slow!

Slow DFT

```
#include <complex>
#include <rarray>
#include <cmath>
using complex = std::complex<double>;

void fft_slow(rvector<complex>& f,
              rvector<complex>& fhat,
              bool inverse)
{
    int n = f.size();
    int sign = inverse?-1:1;
    double v = sign*2*M_PI/n;
    for (int q = 0; q < n; q++)
    {
        fhat[q] = 0.0;
        for (int m = 0; m < n; m++) {
            fhat[q] += complex(cos(v*q*m), sin(v*q*m))
                        * f[m];
        }
    }
}
```

The inverse left out the $1/n$ normalization; this is common in many implementations.

Fast Fourier Transform (FFT)

- Even Gauss realized $\mathcal{O}(n^2)$ was too slow, and came up with a fast version.
- This fast version was derived in partial form several times before and even after Gauss, because he'd just written it in his diary in 1805 (published later).
- Rediscovered (in general form) by Cooley and Tukey in 1965.

Basic idea

- Write each n -point FT as a sum of two $\frac{n}{2}$ point FTs.
- Do this recursively $2 \log n$ times.
- Each level requires $\sim n$ computations: $\mathcal{O}(n \log n)$ instead of $\mathcal{O}(n^2)$.
- Could as easily divide into 3, 5, 7, ... parts. In practice, nobody does.

Fast Fourier Transform: How is it done?

- Define $\omega_n = e^{2\pi i/n}$.
- Note that $\omega_n^2 = \omega_{n/2}$.
- DFT takes form of matrix-vector multiplication:

$$\hat{f}_q = \sum_{j=0}^{n-1} \omega_n^{qj} f_j$$

- With a bit of rewriting (assuming n is even):

$$\hat{f}_q = \underbrace{\sum_{j=0}^{n/2-1} \omega_{n/2}^{qj} f_{2j}}_{\text{FT of even samples}} + \omega_n^q \underbrace{\sum_{j=0}^{n/2-1} \omega_{n/2}^{qj} f_{2j+1}}_{\text{FT of odd samples}}$$

- Repeat, until the lowest level (for $n = 1$, $\hat{f} = f$).
- Note that a fair amount of shuffling is involved.

Inverse DFT

- Inverse DFT is similar to forward DFT, up to a normalization: almost just as fast.

$$f_j = \frac{1}{n} \sum_{q=0}^{n-1} \hat{f}_q e^{\mp 2\pi i j q/n}$$

- FFT allows quick back-and-forth between space and wavenumber domain, or time and frequency domain.
- Allows parts of the computation and/or analysis to be done in the most convenient or efficient domain.

Fast Fourier Transform: Already done!

We've said it before and we'll say it again: Do not write your own: use existing libraries!

Why not write your own?

- Because getting all the pieces right is tricky;
- Getting it to compute fast requires intimate knowledge of how processors work and access memory;
- Because there are libraries available.

Examples:

- ▶ FFTW3 (Faster Fourier Transform in the West, version 3)
 - ▶ cuFFT
 - ▶ Intel MKL
 - ▶ IBM ESSL
-
- Because you have better things to do.

Example of using a library: FFTW

Previous version:

```
#include <complex>
#include <rarray>
#include <cmath>
using complex = std::complex<double>;

void fft_slow(rvector<complex>& f,
              rvector<complex>& fhat,
              bool inverse)
{
    int n = f.size();
    int sign = inverse?-1:1;
    double v = sign*2*M_PI/n;
    for (int q = 0; q < n; q++)
    {
        fhat[q] = 0.0;
        for (int m = 0; m < n; m++) {
            fhat[q] += complex(cos(v*q*m),sin(v*q*m))
                        * f[m];
        }
    }
}
```

FFTW version:

```
#include <complex>
#include <rarray>
#include <fftw3.h>
using complex = std::complex<double>;

void fft_fast(rvector<complex>& f,
              rvector<complex>& fhat,
              bool inverse)
{
    int n = f.size();
    int sign = inverse?FFTW_BACKWARD:FFTW_FORWARD;
    fftw_plan p = fftw_plan_dft_1d(n,
                                   (fftw_complex*)(f.data()),
                                   (fftw_complex*)(fhat.data()),
                                   sign,
                                   FFTW_ESTIMATE);
    fftw_execute(p);
    fftw_destroy_plan(p);
}
```

Notes

- Creates a plan first. This is a mandatory step for `fftw`.
- An `fftw_plan` contains all information necessary to compute the transform, including the pointers to the input and output arrays.
- FFTW uses its own complex number type, completely compatible with C++'s complex numbers, except C++ does not know that. So, casts.
- Plans can be reused in the program, and even saved on disk!
- When creating a plan, you can have FFTW measure the fastest way of computing dft's of that size (FFTW_MEASURE), instead of guessing (FFTW_ESTIMATE).
- Link with `-lfftw3` (for double precision).
- For single precision, use `fftwf_` functions and link with `-lfftw3f`. precision too.

Example

- Create a 1d input signal: a discretized $\text{sinc}(x) = \sin(x)/x$ with 16384 points on the interval $[-30:30]$.
- Perform forward transform
- Write to standard out
- Compile, and linking to fftw3 library.
- Continuous FT of $\text{sinc}(x)$ is the rectangle function:

$$\text{rect}(f) = \begin{cases} 0.5 & \text{if } \|k\| \leq 1 \\ 0 & \text{if } \|k\| > 1 \end{cases}$$

up to a normalization.

- Does it match?

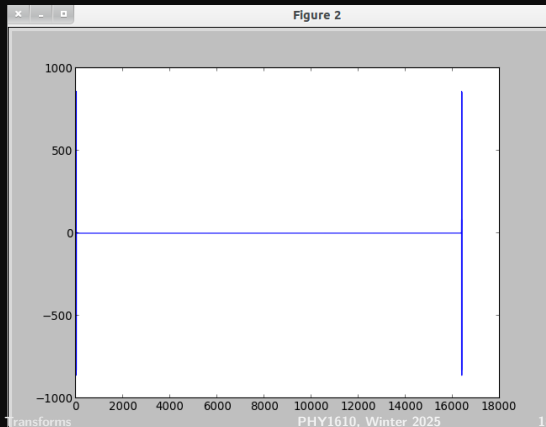
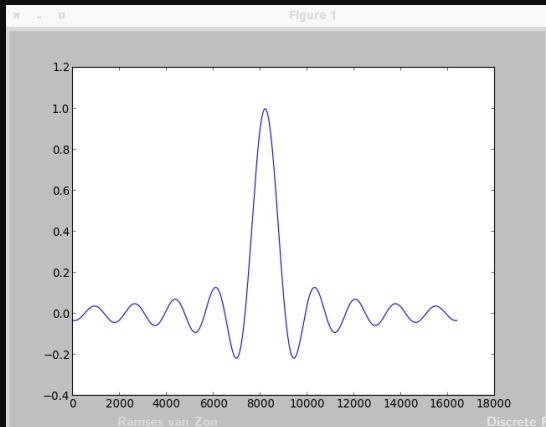
Code for the working example

```
//sincfftw.cpp
#include <iostream>
#include <complex>
#include <rarray>
#include <fftw3.h>
using complex = std::complex<double>;
int main() {
    const int n = 16384;
    rvector<complex> f(n), fhat(n);
    for (int i=0; i<n; i++) {
        double x = 60*(i/double(n)-0.5); // x-range from -30 to 30
        if (x!=0.0) f[i] = sin(x)/x; else f[i] = 1.0;
    }
    fftw_plan p = fftw_plan_dft_1d(n,
                                   reinterpret_cast<fftw_complex*>(f.data()),
                                   reinterpret_cast<fftw_complex*>(fhat.data()),
                                   FFTW_FORWARD, FFTW_ESTIMATE);
    fftw_execute(p);
    fftw_destroy_plan(p);
    for (int i=0; i<n; i++)
        std::cout << f[i].real() << " " << fhat[i].real() << std::endl;
}
```

Compile, link, run, plot

```
$ module load gcc/13 rarray fftw/3 python/3
$ g++ -std=c++17 -c -O3 sincfftw.cpp -o sincfftw.o
$ g++ sincfftw.o -o sincfftw -lfftw3
$ ./sincfftw > output.dat
$ ipython --pylab
```

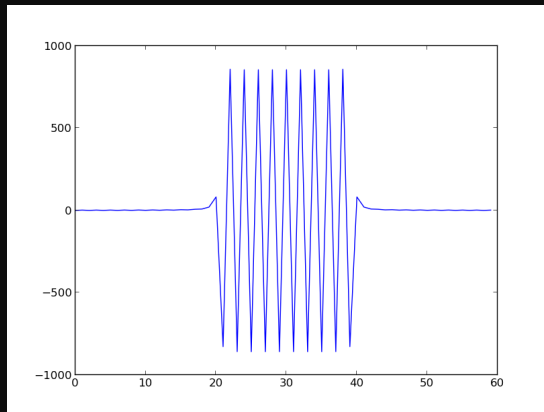
```
>>> data = genfromtxt('output.dat')
>>> plot(data[:,0])
>>> figure()
>>> plot(data[:,1])
```



Plots of the output, rewrapped

Pick the first and the last 30 points.

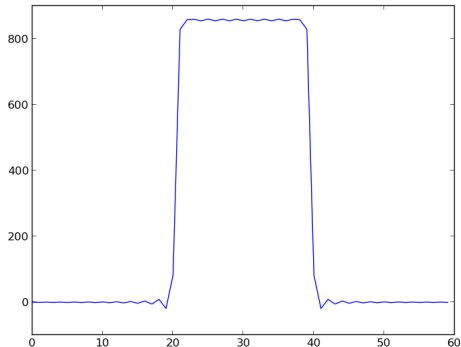
```
>>> x1=range(30)
>>> x2=range(len(data)-30,len(data))
>>> y1=data[x1,1]
>>> y2=data[x2,1]
>>> figure()
>>> plot(hstack((y2,y1)))
```



Undo phase factor due to shifting

```
>>> plot(hstack((y2,y1))*array([1,-1]*30))
```

We retrieved our rectangle function!



Multidimensional transforms

In principle a straightforward generalization:

- Given a set of $n \times m$ function values on a regular grid:

$$f_{ab} = f(a\Delta x, b\Delta y)$$

- Transform these to n other values \hat{f}_{kl}

$$\hat{f}_{kl} = \sum_{a=0}^{n-1} \sum_{b=0}^{m-1} f_{ab} e^{\pm 2\pi i (a k + b l)/n}$$

- Easily back-transformed:

$$f_{ab} = \frac{1}{nm} \sum_{k=0}^{n-1} \sum_{l=0}^{m-1} \hat{f}_{kl} e^{\mp 2\pi i (a k + b l)/n}$$

- Negative frequencies: $f_{-k, -l} = f_{n-k, m-l}$.

Multidimensional FFT

- We could successive apply the FFT to each dimension
- This may require transposes, can be expensive.
- Alternatively, could apply FFT on rectangular patches.
- Mostly should let the libraries deal with this.
- FFT scaling still $n \log n$.

Symmetries for real data

- All arrays were complex so far.
- If input f is real, this can be exploited.

$$f_j^* = f_j \leftrightarrow \hat{f}_k = \hat{f}_{n-k}^*$$

- Each complex number holds two real numbers, but for the input f we only need n real numbers.
- If n is even, the transform \hat{f} has real \hat{f}_0 and $\hat{f}_{n/2}$, and the values of $\hat{f}_k > n/2$ can be derived from the complex valued $\hat{f}_{0 < k < n/2}$: again n real numbers need to be stored.

Symmetries for real data

- A different way of storing the result is in “half-complex storage’’. First, the $n/2$ real parts of $\hat{f}_{0 \leq k < n/2}$ are stored, then their imaginary parts in reversed order.
- Seems odd, but means that the magnitude of the wave-numbers is like that for a complex-to-complex transform.
- These kind of implementation dependent storage patterns can be tricky, especially in higher dimensions.

Applications

Application of the Fourier transform

- ① Signal processing, certainly.
- ② Many equations become simpler in the fourier basis.
 - ▶ Reason: $\exp(ikx)$ are eigenfunctions of the $\partial/\partial x$ operator.
 - ▶ Partial diferential equation become algebraic ones, or ODEs.
 - ▶ Thus avoids matrix operations.
- ③ Optimizing long range particle-particle interactions in N-body simulations and molecular dynamics.

2. Solving diffusion equation with FFT

$$\frac{\partial u}{\partial t} = \kappa \frac{\partial^2 u}{\partial x^2}$$

for $u(x, t)$ on $x \in [0, L]$, with boundary conditions $u(0, t) = u(L, t) = 0$, and $u(x, 0) = f(x)$.

Write

$$u(x, t) = \sum_{k=-\infty}^{\infty} \hat{u}_k(t) e^{2\pi i k x / L}$$

then the PDE becomes an ODE:

$$\frac{d\hat{u}_k}{dt} = -\kappa \frac{4\pi^2 k^2}{L^2} \hat{u}_k; \quad \text{with } \hat{u}_k(0) = \hat{f}_k.$$

Alternatively, one can first discretize the PDE, then take an FFT. This is numerically different.

3. Long-range particle interactions

- Long-range interactions are those that cannot be cut off without seriously altering the physics.
- Examples of a long range interactions include:
 - ▶ Gravity
 - ▶ Electrostatics
- In N-body and MD simulations, the force computation is often the bottleneck.
- Without a cut-off (as for short-range) interactions, we are left with a sum over interacting pairs, i.e., an or “Particle-Particle”, $\mathcal{O}(N^2)$ method.
- As we saw in the Molecular Dynamics lecture, we can use Particle-Mesh and solve part of the interactions in fourier space; $\mathcal{O}(N \log N)$